

9 Tall buildings

9.1 Introduction

Tall buildings, now approaching 500 m in height, project well into the atmospheric boundary layer, and their upper levels may experience the highest winds of large-scale windstorms, such as tropical cyclones or the winter gales of the temperate regions. Resonant dynamic response in along-wind, cross-wind and torsional modes are a feature of the overall structural loads experienced by these structures. Extreme local cladding pressures may be experienced on their side walls.

The post World War II generation of high-rise buildings were the stimulus for the development of the boundary-layer wind tunnel, which remains the most important tool for the establishment of design wind loads on major building projects in many countries.

In this chapter, the history of investigations into wind loading of tall buildings, the major response mechanisms and phenomena, and the available analytical and semi-analytical techniques, will be discussed.

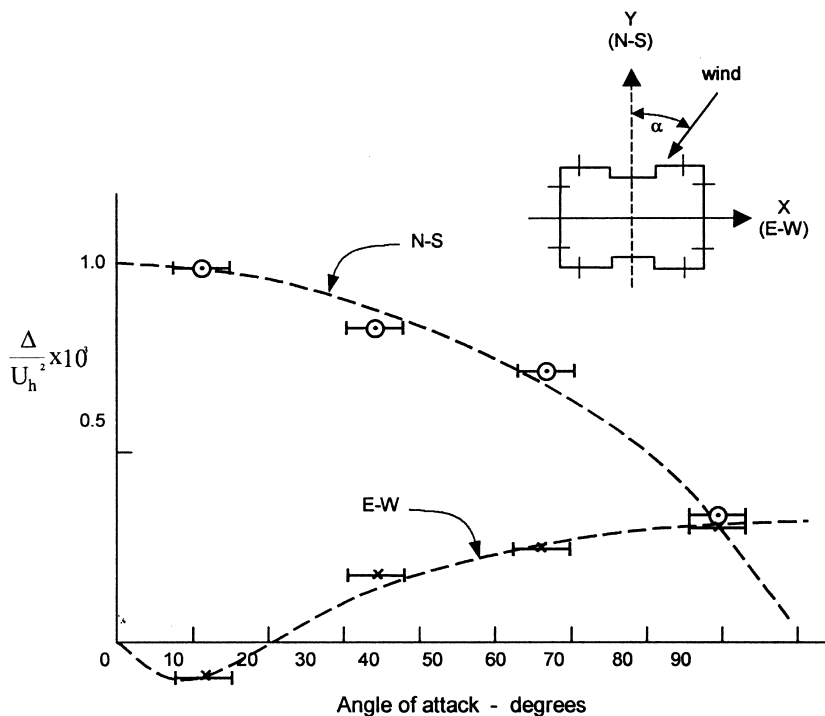
9.2 Historical

Tall buildings, or ‘skyscrapers’ are amongst the more wind-sensitive of structures, and it was inevitable that their response to wind would be of concern to structural engineers, and attract the interest of early experimenters, both in the wind tunnel, and in full scale.

The Empire State Building, at 380 m, was the tallest building in the world for forty years, and was the subject of three significant studies in the 1930s (Coyle, 1931; Dryden and Hill, 1933; Rathbun, 1940). These studies have been re-appraised in some detail by Davenport (1975).

Coyle (1931) used a portable horizontal pendulum to record the motion of the building. This clearly revealed resonant dynamic response with a period of around 8 seconds. Rathbun’s (1940) extensive full scale measurements were described by Davenport as: ‘a monumental piece of full-scale experimentation’. Wind pressures on three floors of the building were measured with 30 manometers and 28 flash cameras. The pressure coefficients showed considerable scatter, but were clearly much lower than those obtained by Dryden and Hill (1933) on a wind tunnel model in a uniform flow some years earlier. Rathbun also performed deflection measurements on the Empire State Building using a plumb bob extending from the 86th floor to the 6th floor. These results (as re-analysed by Davenport) indicated the significantly different stiffness of the building in the east–west direction in comparison with the north–south direction ([Figure 9.1](#)).

In the 1960s and 1970s, a resurgence in the building of skyscrapers occurred – particularly in North America, Japan and Australia. There was great interest in wind loads on tall buildings at this time – this has continued to the end of the twentieth century. The two main problem areas to emerge were:



Δ - Mean deflection (inches)

U_h - Mean wind speed at 1250 feet MPH (uncorrected)

Figure 9.1 Full-scale measurements of mean deflection on the Empire State Building by Rathbun (1940) – reanalysed by Davenport (1975).

- The vulnerability of glazed cladding to both direct wind pressures, and flying debris in windstorms
- Serviceability problems arising from excessive motion near the top of tall buildings

From the early 1970s, many new building proposals were tested in the new boundary-layer wind tunnels (see [Chapter 7](#)), and quite a few full-scale monitoring programmes were commenced.

One of the most comprehensive and well-documented full-scale measurement studies, with several aspects to it, which lasted for most of the 1970s, was that on the 239 m tall Commerce Court building in Toronto, Canada (Dalglish, 1975; Dalglish *et al.*, 1979; Dalglish *et al.*, 1983). The full scale studies were supplemented with wind tunnel studies, both in the design stage (Davenport *et al.*, 1969) and later on a pressure model (Dalglish *et al.*, 1979), and a multi-degree-of-freedom aeroelastic model, in parallel with the full scale studies (Templin and Cooper, 1981; Dalglish *et al.*, 1983).

The early full-scale pressure measurements on the Commerce Court building showed good agreement with the wind tunnel study (at 1/400 scale) for mean pressure coefficients, and for the mean base shear and overturning moment coefficients. Not as good agreement with the 1/400 scale wind tunnel tests, was found for the r.m.s. fluctuating pressure coef-

ficients for some wind directions (Dalglish, 1975). The later reported pressure measurements (Dalglish *et al.*, 1979) showed better agreement for the fluctuating pressure and peak measurements on a larger (1/200) scale wind tunnel model, with accurately calibrated tubing and pressure measurement system. The full-scale pressure study on Commerce Court highlighted the importance of short duration peak pressures in separated flow regions (at around this time similar observations were being made from the roof of the low-rise building at Aylesbury – Section 8.2.2). Subsequently, detailed statistical studies of these were carried out for application to glass loading (see Section 9.4.5). Although the Commerce Court pressure measurements were of a high quality, they suffered from the lack of an independent reference pressure for the pressure coefficients – an internal pressure reading from the building was used. For comparison of mean pressure coefficients with the wind tunnel results, it was necessary to force agreement at one pressure tapping – usually in wake region.

The full-scale study of acceleration response (Dalglish *et al.*, 1983) showed the following features:

- the significance of the torsional (twisting) motions superimposed on the sway motions for one direction (E–W). This was explained by an eccentricity in the north–south direction between the centre of mass, and the elastic axis
- generally good agreement between the final aeroelastic model, which included torsional motions, and the full scale data, for winds from a range of directions
- reasonable agreement between the full-scale data and predictions of the National Building Code of Canada for along – and cross-wind accelerations.

The agreements observed occurred despite some uncertainties in the reference velocity measured at the top of the building, and in the dynamic properties (frequency and damping) of the building. An interesting observation, not yet clearly explained, but probably an added mass effect, was a clear decrease in observed building frequency as the mean speed increased.

Another important full-scale study, significant for its influence on the development of the British Code of Practice for Wind Loads, was that carried out on the 18-storey Royex House in London (Newberry *et al.*, 1967). This study revealed aspects of the transient and fluctuating pressures on the windward and side walls.

The first major boundary-layer wind tunnel study of a tall building was that carried out for the twin towers of the World Trade Center, New York, in the mid 1960s, at Colorado State University. This was the first of many commercial studies, now numbering in the thousands, in boundary-layer wind tunnels.

9.3 Flow around tall buildings

Tall buildings are bluff bodies of medium to high aspect ratio, and the basic characteristics of flow around this type of body were covered in some detail in [Chapter 4](#). [Figure 9.2](#) shows the general characteristics of boundary-layer wind flow around a tall building. On the windward face there is a strong downward flow below the stagnation point, which occurs at a height of 70 to 80% of the overall building height. The down flow can often cause problems at the base, as high velocity air from upper levels is brought down to street level. Separation and re-attachment at the side walls are associated with high local pressures. The rear face is a negative pressure region of lower magnitude mean pressures, and a low level of fluctuating pressures.

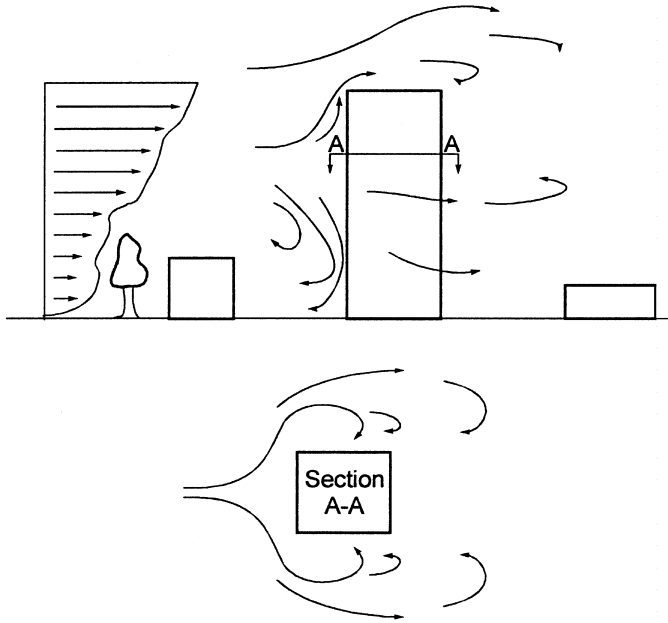


Figure 9.2 Wind flow around a tall building.

In a mixed extreme wind climate of thunderstorm downbursts (Section 1.3.5) and synoptic winds, the dominant wind for wind loading of tall buildings will normally be the latter, as the downburst profile has a maximum at a height of 50–100 m (Figure 3.3).

9.4 Cladding pressures

9.4.1 Pressure coefficients

As in previous chapters, pressure coefficients in this chapter will be defined with respect to a mean wind speed at the top of the building, denoted by \bar{U}_h . Thus, the mean, root-mean-square fluctuating (standard deviation), maximum and minimum pressure coefficients are defined according to equations (9.1), (9.2), (9.3) and (9.4), respectively.

$$\bar{C}_p = \frac{\bar{p} - p_0}{\frac{1}{2}\rho_a \bar{U}_h^2} \quad (9.1)$$

$$C'_p = \sigma_{C_p} = \frac{\sqrt{\bar{p'^2}}}{\frac{1}{2}\rho_a \bar{U}_h^2} \quad (9.2)$$

$$\hat{C}_p = \frac{\hat{p} - p_0}{\frac{1}{2}\rho_a \bar{U}_h^2} \quad (9.3)$$

$$\check{C}_p = \frac{\check{p} - p_0}{\frac{1}{2}\rho_a \bar{U}_h^2} \quad (9.4)$$

In equations (9.3) and (9.4), the maximum and minimum pressures, \hat{p} and \check{p} , are normally defined as the average or expected peak pressure at a point in a given averaging time, which may be taken as a period between 10 minutes and 3 hours in full scale. It is not usually convenient, or economic, to measure such average peaks directly in wind tunnel tests, and various alternative statistical procedures have been proposed. These are discussed in Section 9.4.4.

9.4.2 Pressure distributions on buildings of rectangular cross-section

The local pressures on the wall of a tall building can be used directly for the design of cladding, which is generally supported over small tributary areas.

Figure 4.15 shows the distribution of mean pressure coefficient on the faces of tall prismatic shape, representative of a very tall building, with aspect ratio (height/width) of 8, in a boundary-layer flow.

Figures 9.3, 9.4 and 9.5 show the variation in mean, maximum and minimum pressure coefficients on the windward, side and leeward faces, for a lower building of square cross-section, with aspect ratio equal to 2.1 (Cheung, 1984). The pressures were measured on a wind tunnel models which represented a building of 85 m height; the building is isolated, that is there is no shielding from buildings of comparable height, and the approaching flow was boundary-layer flow over suburban terrain. The value of Jensen number, h/z_0 , (see Section 4.4.4) was then approximately 40.

Figure 9.3 shows a stagnation point on the windward face, where the value of \bar{C}_p reaches

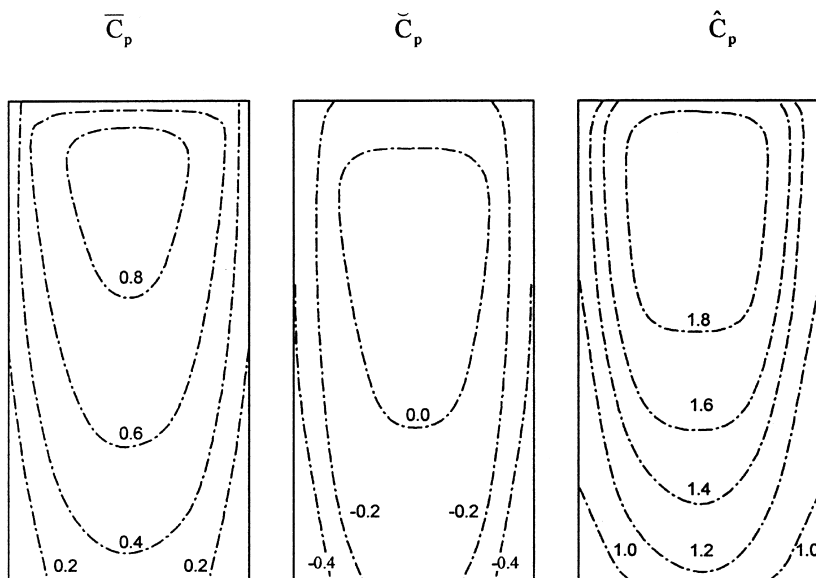


Figure 9.3 Mean, maximum and minimum pressure coefficients – windward wall of a building with square cross section – height/width = 2.1 (Cheung, 1984).

a maximum, at about 0.8 h. The heights for largest maximum pressure coefficient are slightly lower than this.

The side walls (Figure 9.4) are adjacent to a flow which is separating from the front wall, and generating strong vortices (see Figures 4.1 and 9.2). The mean pressure coefficients are generally in the range from -0.6 to -0.8 , and not dissimilar to the values on the much taller building in Figure 4.15. The largest magnitude minimum pressure coefficients of about -3.8 occur near the base of the buildings, unlike the windward wall pressures. A wind direction parallel to the side wall produces the largest magnitude negative pressures in this case.

The mean and largest peak pressures on the leeward wall (Figure 9.5) are also negative, but are typically half the magnitude of the side wall pressures. This wall is of course sheltered, and exposed to relatively slowly moving air in the near wake of the building.

9.4.3 The nature of fluctuating local pressures and probability distributions

As discussed in Section 9.2, in the 1970s, full-scale and wind tunnel measurements of wind pressures on tall buildings, highlighted the local peak negative pressures, that can occur, for some wind directions, on the walls of tall buildings, particularly on side walls at locations near windward corners, and on leeward walls. These high pressures generally only occur for quite short periods of time, and may be very intermittent in nature. An example of the intermittent nature of these pressure fluctuations is shown in Figure 9.6 (from Dalgleish, 1971).

Several studies (e.g. Dalgleish, 1971; Peterka and Cermak, 1975) indicated that the probability densities of pressure fluctuations in separated flow regions on tall buildings were not well fitted by the normal or Gaussian probability distribution (Appendix E). This is the case, even though the latter is a good fit to the turbulent velocity fluctuations in the

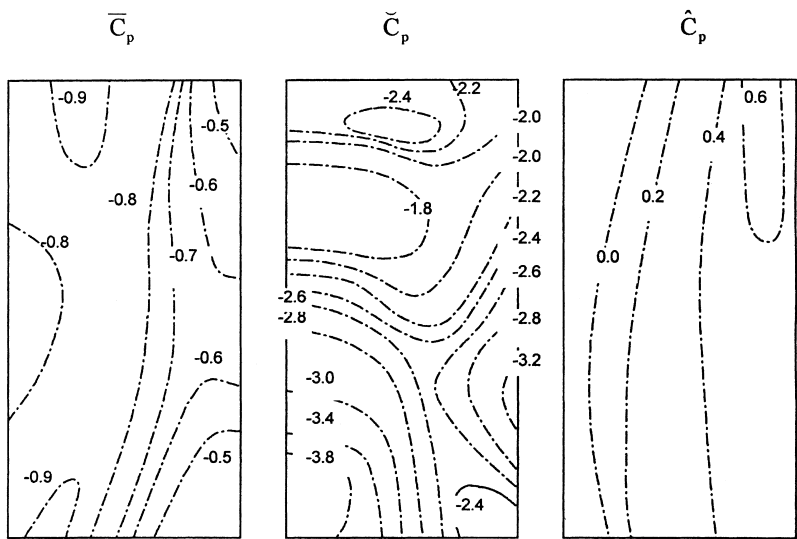


Figure 9.4 Mean, maximum and minimum pressure coefficients – side wall of a building with square cross section – height/width = 2.1 (Cheung, 1984).

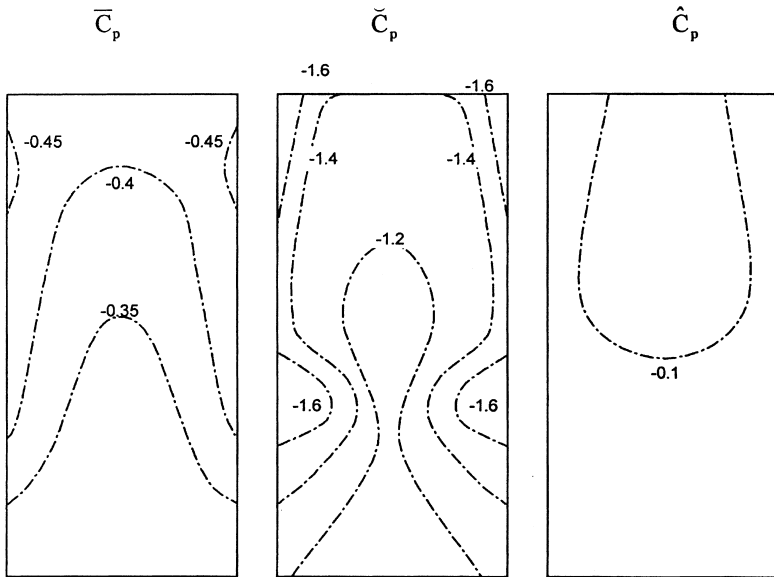


Figure 9.5 Mean, maximum and minimum pressure coefficients – leeward wall of a building with square cross section – height/width = 2.1 (Cheung, 1984).

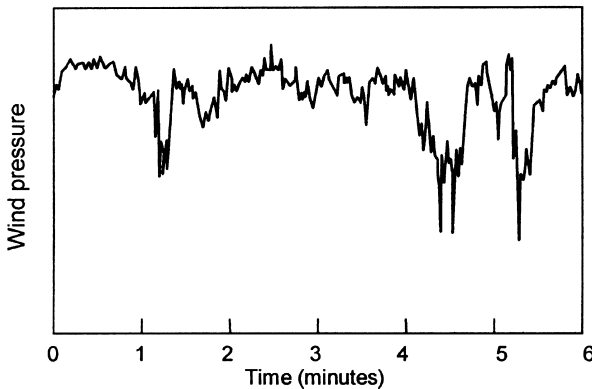


Figure 9.6 Record of fluctuating pressure from the leeward wall of a full-scale office building (Dalglish, 1971).

wind (see Section 3.3.2). The ‘spiky’ nature of local pressure fluctuations (Figure 9.6) results in probability densities of peaks of five standard deviations, or greater, below the mean pressure, being several times greater than that predicted by the Gaussian distribution. This is illustrated in Figure 9.7 derived from wind tunnel tests of two tall buildings (Peterka and Cermak, 1975).

A consequence of the intermittency and non-Gaussian nature of pressure fluctuations on tall buildings, is that the maximum pressure coefficient measured at a particular location on a building in a defined time period – say 10 minutes in full scale – may vary consider-

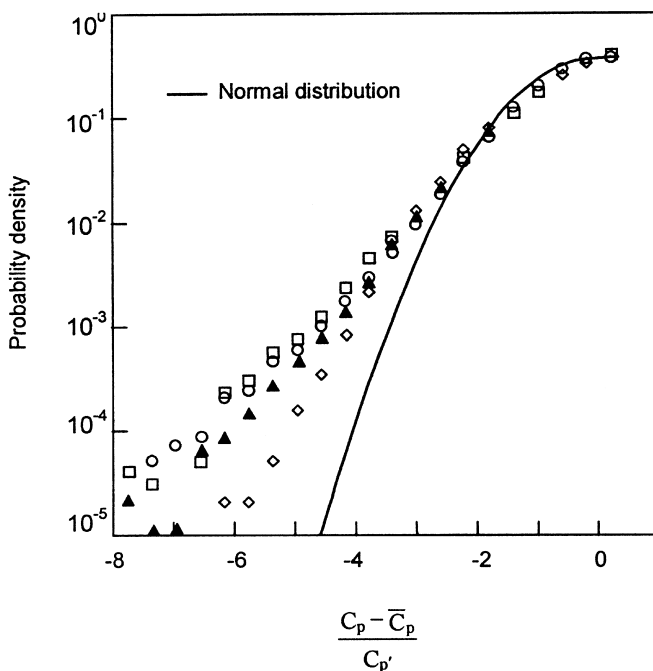


Figure 9.7 Probability densities of pressure fluctuations from regions in separated flow on tall buildings (Peterka and Cermak, 1975).

ably from one time period to the next. Therefore they cannot be predicted by knowing the mean and standard deviation, as is the case with a Gaussian random process (Davenport, 1964). This has led to a number of different statistical techniques being adopted to produce more consistent definitions of peak pressures for design – these are discussed in Section 9.4.4. A related matter is the response characteristics of glass cladding to short duration peak loads. The latter aspect is discussed in Section 9.4.5.

A detailed study (Surry and Djakovich, 1995) of local negative peak pressures on generic tall building models of constant cross-section, with four different corner geometries, indicated that the details of the corner geometry do not affect the general magnitude of the minimum pressure coefficients, but rather the wind direction at which they occur. The highest peaks were associated with vortex shedding.

9.4.4 Statistical methods for determination of peak local pressures

A simple approach, originally proposed by Lawson (1976), uses the parent probability distribution of the pressure fluctuations, from which a pressure coefficient, with a designated (low) probability of exceedence is extracted. The probability of exceedence is normally in the range 1×10^{-4} to 5×10^{-4} , with the latter being suggested by Lawson. This method can be programmed ‘on the run’ in wind tunnel tests, relatively easily; sometimes a standard probability distribution, such as the Weibull type (see [Appendix C3.4](#)) is used to fit the measured data and interpolate, or extrapolate, to the desired probability level.

Cook and Mayne (1979) proposed a method in which the total averaging time, T , is

divided into sixteen equal parts and the measured peak pressure coefficient (maximum or minimum) within each reduced time period, t , is retained. A Type I Extreme Value (Gumbel) distribution (see Section 2.2.1 and [Appendix C4](#)) is fitted to the measured data, giving a mode, c_t , and scale factor, a_t . These can then be used to calculate the parameters of the Extreme Value Type I distribution appropriate to the maxima (or minima) for the original time period, T , as follows:

$$c_T = c_t + a_T \log_e 16 \quad (9.5)$$

$$a_T = a_t \quad (9.6)$$

Knowing the distribution of the extreme pressure coefficients, the expected peak, or any other percentile, can then be easily determined. The method proposed by Cook and Mayne (1979), in fact, proposes an effective peak pressure coefficient C_p^* given by:

$$C_p^* = c_T + 1.4a_T \quad (9.7)$$

Peterka (1983) proposed the use of the probability distribution of 100 independent maxima within a time period equivalent to 1 h, to determine C_p^* .

Another approach is to make use of level crossing statistics. Melbourne (1977) proposed the use of a normalised rate of crossing of levels of pressure (or structural response). A nominal rate of crossing (e.g. 10^{-4} per hour) is chosen to determine a nominal level of 'peak' pressure.

The parameters of the (Type I) extreme value distribution for the extreme pressure in a given time period can also be derived from level crossing rates as follows. The level crossings are assumed to be uncorrelated events which can be modelled by a Poisson distribution ([Appendix C3.5](#)).

The Poisson distribution gives the probability for the number of events, n , in a given time period, T , when the average rate of occurrence of the events is v :

$$P(n,v) = \frac{(vT)^n}{n!} \exp(-vT) \quad (9.8)$$

The 'event' in this case can be taken as an upcrossing of a particular level, e.g. the exceedence of a particular pressure level. The probability of getting no crossings of a pressure level, p , during the time period, T , is also the probability that the largest value of the process $p(t)$, during the time period, is less than that level, i.e. the cumulative probability distribution of the largest value in the time period, T .

Thus,

$$F(p) = P(0,v) = \frac{(vT)^0}{0!} \exp(-vT) = \exp(-vT) \quad (9.9)$$

If we assume that the average number of crossings of level x in time T , is given by:

$$vT = \exp\left[-\frac{1}{a}(p-u)\right] \quad (9.10)$$

where a and u are constants, then,

$$F(x) = \exp\left\{\exp\left[-\frac{1}{a}(p - u)\right]\right\} \quad (9.11)$$

This is the Type I (Gumbel) extreme value distribution with a mode of u and a scale factor of a .

From equation (9.10), taking natural logarithms of both sides,

$$\log_e(vT) = -\frac{1}{a}(p - u) \quad (9.12)$$

The mode and scale factor of the Type I extreme value distribution of the process $p(t)$ can be estimated by the following procedure:

- Plot the natural logarithm of the rate of upcrossings against the level, p
- Fit a straight line. From equation (9.12), the slope is $(-1/a)$, and the intercept ($p = 0$) is (u/a)
- From these values, estimate u and a , the mode and scale factor of the Type I extreme value distribution of p .

9.4.5 Strength characteristics of glass in relation to wind loads

Direct wind loading is a major design consideration in the design of glass and its fixing in tall buildings. However, the need to design for wind-generated flying debris (Section 1.5) – particularly roof gravel – in some cities, also needs to be considered (Minor, 1994).

As has been discussed, wind pressures on the surfaces of buildings fluctuate greatly with time, and it is known that the strength of glass is quite dependent on the duration of the loading. The interaction of these two phenomena results in a complex design problem.

The surfaces of glass panels are covered with flaws of various sizes and orientations. When these are exposed to tensile stresses they grow at a rate dependent on the magnitude of the stress field, as well as relative humidity and temperature. The result is a strength reduction which is dependent on the magnitude and duration of the tensile stress. Drawing on earlier studies of this phenomenon, known as ‘static fatigue’, Brown (1972) proposed a formula for damage accumulation which has the form of equation (9.13), at constant humidity and temperature.

$$D = \int_0^T [s(t)]^n dt \quad (9.13)$$

where D is the accumulated damage, $s(t)$ is the time varying stress, T is the time over which the glass is stressed, n is a higher power (in the range of 12 to 20).

The expected damage, in time T , under a fluctuating wind pressure $p(t)$, in the vicinity of a critical flaw can be written as equation (9.14).

$$E\{D\} = K \int_0^T E\{[p(t)]^m\} dt \quad (9.14)$$

where K is a constant, and m is a different power, usually lower than n , but dependent on the size and aspect ratio of the glass, which allows for the non-linear relationship between load and stress for glass plates due to membrane stresses (Calderone and Melbourne, 1993). $E\{\}$ is the expectation or averaging operation.

Calderone (1999), after extensive glass tests, found a power law relationship between maximum stress anywhere in a plate, and the applied pressure, for any given plate; this may be used to determine the value of m for that plate. Values fall in the range of 5 to 20.

The integral on the right-hand-side of equation (9.14) is T times the m th moment of the pressure fluctuation, so that:

$$E\{D\} = KT \left(\frac{1}{2} \rho \bar{U}^2 \right)^m \int_0^{\infty} C_p^m f_{C_p}(C_p) dC_p \quad (9.15)$$

where $C_p(t)$ is the time-varying pressure coefficient, and $f_{C_p}(C_p)$ is the probability density function for C_p .

The integral in equation (9.15) is proportional to the rate at which damage is accumulated in the glass panel. It can be evaluated from known or expected probability distributions (e.g. Holmes, 1985), or directly from wind tunnel or full-scale pressure-time histories (Calderone and Melbourne, 1993).

The high weighting given to the pressure coefficient by the power, m , in equation (9.15) means that the main contribution to glass damage comes from isolated peak pressures, which typically occur intermittently on the walls of tall buildings (see Figure 9.6).

An equivalent static pressure coefficient, C_{ps} , which corresponds to a constant pressure which gives the same rate of damage accumulation as a fluctuating pressure-time history, can be defined:

$$C_{ps} = \left[\int_0^{\infty} C_p^m f_{C_p}(C_p) dC_p \right]^{1/m} \quad (9.16)$$

For the structural design of glazing, it is necessary to relate the computed damage caused by wind action, to failure loads obtained in laboratory tests of glass panels. The damage integral (equations (9.13) or (9.14)), can be used to compute the damage sustained by a glass panel under the 'ramp' loading (i.e. increasing linearly with time) commonly used in laboratory testing. In these tests, failure typically occurs in about 1 min.

An equivalent glass design coefficient, C_k , can be defined (Dalglish, 1979) which, when multiplied by the reference dynamic pressure, $(\frac{1}{2} \rho_a \bar{U}^2)$, gives a pressure which produces the same damage in a 60 second ramp increase, as in a windstorm of specified duration.

Making use of equations (9.15) and (9.16), it can be easily shown that for a windstorm of 1 h duration:

$$C_k = [60(1 + m)]^{1/m} C_{ps} \quad (9.17)$$

Using typical values of m and typical probability distributions, it can be shown (Dalglish, 1979; Holmes, 1985) that C_k is approximately equal to the expected peak pressure coefficient occurring during the hour of storm wind. This fortuitous result, which is insensitive to both the value of m and the probability distribution, means that measured peak pressure

coefficients from wind tunnel tests are valid for use in calculation of design loads, for comparison with 1-min loads in glass design charts.

9.5 Overall loading and dynamic response

In [Chapter 6](#), the random or spectral approach to the along-wind response of tall structures was discussed. This approach is widely used for the prediction of the response of tall office buildings in simplified forms in codes and standards (see [Chapter 15](#)). Dynamic response of a tall building in the along-wind direction is primarily produced by the turbulent velocity fluctuations in the natural wind (Section 3.3). In the cross-wind direction, loading and dynamic response is generated by random vortex shedding (Section 4.6.3) – that is, it is a result of unsteady separating flow generated by the building itself, with a smaller contribution from cross-wind turbulence.

9.5.1 General response characteristics

In this section some general characteristics of the dynamic response of tall buildings to wind will be outlined.

By a dimensional analysis, or by application of the theory given in Section 5.3.1, it can be demonstrated (Davenport, 1966, 1971) that the root-mean-square fluctuating deflection at the top of a tall building of given geometry in a stationary (synoptic) wind, is given to a good approximation for the along-wind response by:

$$\frac{\sigma_x}{h} = A_x \left(\frac{\rho_a}{\rho_b} \right) \left(\frac{\bar{U}_h}{n_1 b} \right)^{k_x} \frac{1}{\sqrt{\eta}} \quad (9.18)$$

and for the cross-wind response:

$$\frac{\sigma_y}{h} = A_y \left(\frac{\rho_a}{\rho_b} \right) \left(\frac{\bar{U}_h}{n_1 b} \right)^{k_y} \frac{1}{\sqrt{\eta}} \quad (9.19)$$

where, h is the building height, A_x , A_y are constants for a particular building shape, ρ_a is the density of air, ρ_b is an average building density, \bar{U}_h is the mean wind speed at the top of the building, b is the building breadth, k_x , k_y are exponents, n_1 is the first mode natural frequency, and η is the critical damping ratio in the first mode of vibration.

Equations (9.18) and (9.19) are based on the assumption that the responses are dominated by the resonant components. For along-wind response, the background component is independent of the natural frequency. In the case of the cross-wind response, there is no mean component, but some background contribution due to cross-wind turbulence. The assumption of dominance of resonance is valid for slender tall buildings with first mode natural frequencies less than about 0.5 Hz, and damping ratios less than about 0.02.

The equations illustrate that the fluctuating building deflection can be reduced by either increasing the building density or the damping. The damping term, η , includes aerodynamic damping as well as structural damping; however this is normally small for tall buildings.

The term $(\bar{U}_h/n_1 b)$ is a non-dimensional mean wind speed, known as the *reduced velocity*. The exponent, k_x , for the fluctuating along-wind deflection is greater than 2, since the spectral density of the wind speed near the natural frequency, n_1 , increases at a greater

power than 2, as does the aerodynamic admittance function (Section 5.3.1 and Figure 5.4) at that frequency. The exponent for cross-wind deflection, k_y , is typically about 3, but can be as high as 4.

Figure 9.8 shows the variation of (σ_x/h) and (σ_y/h) with reduced velocity for a building of circular cross section (as well as the variation of \bar{X}).

9.5.2 Effect of building cross-section

In a study used to develop an optimum building shape for the U.S. Steel building, Pittsburgh, the response of six buildings of identical height and dynamic properties, but with different cross-sections were investigated in a boundary-layer wind tunnel (Davenport, 1971). The probability distributions of the extreme responses in a typical synoptic wind climate was determined, and are shown plotted in Figure 9.9. The figure shows a range of 3:1 in the responses with a circular cross-section producing the lowest response, and an equilateral triangular cross-section the highest. Deflection across the shortest (weakest) axis of a 2:1 rectangular cross-section was also large.

9.5.3 Corner modifications

Slotted and chamfered corners on rectangular building cross-sections have significant effects on both along-wind and cross-wind dynamic responses to wind (Kwok and Bailey, 1987; Kwok *et al.*, 1988; Kwok, 1995). Chamfers of the order of 10% of the building

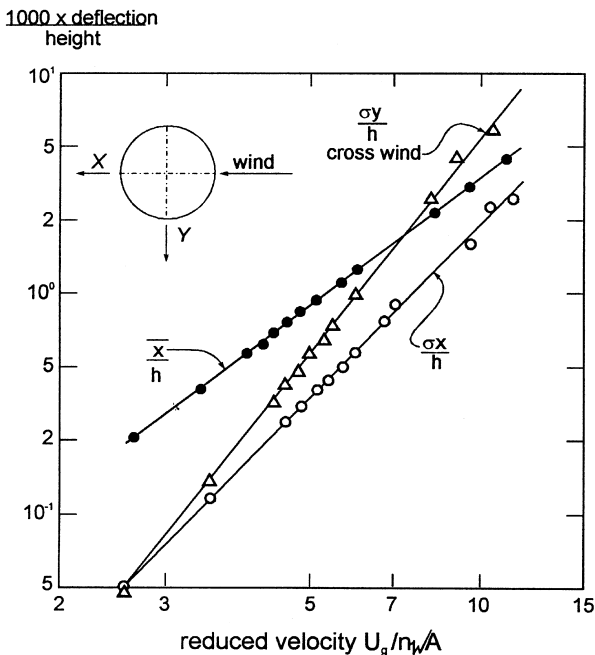


Figure 9.8 The mean and fluctuating response of a tall building of circular cross-section (from Davenport, 1971).

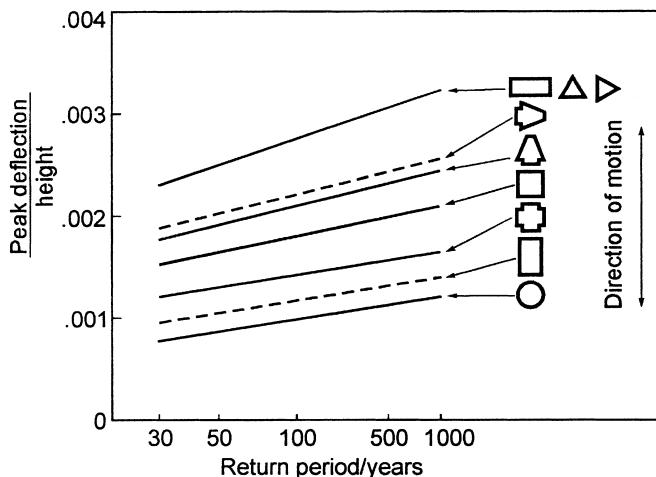


Figure 9.9 Effect of cross-sectional shape on maximum deflections of six buildings (Davenport, 1971).

width produce up to 40% reduction in the along-wind response and 30% reduction in the cross-wind response.

9.5.4 Prediction of cross-wind response

Along-wind response of isolated tall buildings can be predicted reasonably well from the turbulence properties in the approaching flow by applying the random vibration theory methods discussed in Section 5.3.1. Cross-wind response however is more difficult to predict, since vortex shedding plays a dominant role in the exciting forces in the cross-wind direction. However, an approach which has been quite successful, is the use of the high-frequency base balance technique to measure the spectral density of the generalised force in wind tunnel tests (Section 7.6.2). Multiplication by the mechanical admittance and integration over frequency can then be performed to predict the building response.

Examples of generalized force spectra for buildings of square cross-section are shown in Figure 9.10 (Saunders, 1974). Non-dimensional spectra for three different height/breadth ratios are shown, and the approach flow is typical of suburban terrain. The mode shapes are assumed to be linear with height. The abscissa of this graph is reduced frequency – the reciprocal of reduced velocity.

For reduced velocities of practical importance (2 to 8), the non-dimensional spectra vary with reduced velocity to a power of 3 to 5, or with reduced frequency to a power of –3 to –5 (represented by the slope on the log-log plot). Such data have been incorporated in the some standards and codes for design purposes (see Section 15.9).

9.6 Combination of along- and cross-wind response

When dealing with the response of tall buildings to wind loading, the question arises: how should the responses in the along- and cross-wind directions be combined statistically? Since clearly the along-wind and cross-wind responses are occurring simultaneously on a structure it would be unconservative (and potentially dangerous!) to treat these as separate

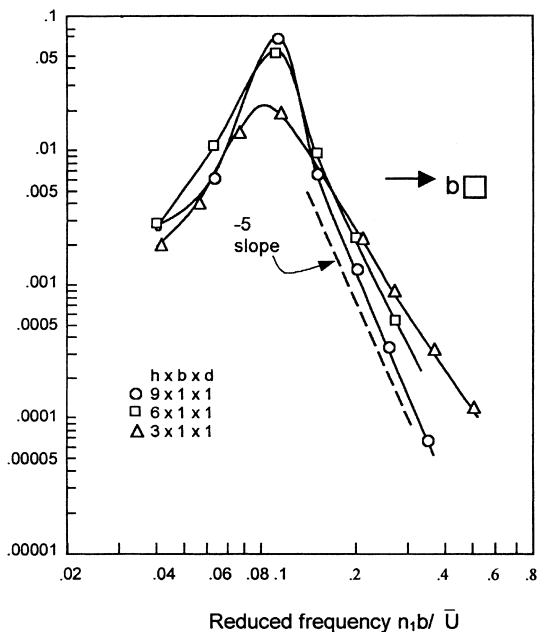


Figure 9.10 Cross-wind generalized force spectra for buildings of square cross-section (Saunders, 1974).

load cases. The question arises when applying those wind loading codes and standards which provide methods for calculating both along-wind and cross-wind dynamic response for tall buildings (see Chapter 15). It also arises when wind tunnel tests are carried out using either aeroelastic (Section 7.6.1), or base-balance methods (Section 7.6.2). In these cases, predictions are usually provided for each wind direction, with respect to *body-* or *building-* axes rather than *wind axes* (see Section 4.2.2. and Figure 4.2). These axes are usually the two principal axes for sway of the building.

Two cases can be identified:

- ‘scalar’ combination rules for load effects
- ‘vector’ combination of responses

The former case is the more relevant case for structural load effects being designed for strength, as in most cases structural elements will ‘feel’ internal forces and stresses from both response directions, and will be developed in the following. The second case is relevant when axi-symmetric structures are under consideration, i.e. structures of circular cross-section such as chimneys.

Load effects (i.e. member forces and internal stresses) resulting from overall building response in two orthogonal directions (*x-* and *y-*) can very accurately be combined by the following formula:

$$\hat{\epsilon}_t = \bar{\epsilon}_x + \bar{\epsilon}_y + \sqrt{(\bar{\epsilon}_x - |\bar{\epsilon}_x|)^2 + (\bar{\epsilon}_y - |\bar{\epsilon}_y|)^2} \quad (9.20)$$

where $\hat{\epsilon}_t$ is total combined maximum peak load effect (e.g. the axial load in a column),

$\bar{\epsilon}_x$ is the load effect derived from the mean response in the x -direction (usually derived from the mean base bending moment in that direction), $\bar{\epsilon}_y$ is the load effect derived from the mean response in the y -direction, $\hat{\epsilon}_x$ is the peak load effect derived from the response in the x -direction and $\hat{\epsilon}_y$ is the peak load effect derived from the response in the y -direction.

Equation (9.20) is quite an accurate one, as it is based on the combination of uncorrelated Gaussian random processes, for which it is exact. Most responses dominated by resonant contributions to wind, have been found to be very close to Gaussian, and if the two orthogonal sway frequencies are well separated, the dynamic responses will be poorly correlated.

As an alternative approximation, the following load cases can be studied:

- (a) [Mean x -load + 0.75(peak – mean) $_x$] with [mean y -load] + 0.75(peak – mean) $_y$]
- (b) [Mean x -load + (peak – mean) $_x$] with [mean y -load]
- (c) [Mean x -load] with [mean y -load + (peak – mean) $_y$]

The case (a) corresponds to the following approximation to equation (9.20) for peak load effect:

$$\epsilon_t = \bar{\epsilon}_x + \bar{\epsilon}_y + 0.75((\hat{\epsilon}_x - |\bar{\epsilon}_x|) + (\hat{\epsilon}_y - |\bar{\epsilon}_y|)) \quad (9.21)$$

Equation (9.21) is a good approximation to equation (9.20) for the range:

$$1/3 < (\hat{\epsilon}_x - |\bar{\epsilon}_x|)/(\hat{\epsilon}_y - |\bar{\epsilon}_y|) < 3$$

The other two cases (b) and (c) are intended to cover the cases outside this range, i.e. when $(\hat{\epsilon}_x - |\bar{\epsilon}_x|)$ is much larger than $(\hat{\epsilon}_y - |\bar{\epsilon}_y|)$, and vice-versa.

9.7 Torsional loading and response

The significance of torsional components in the dynamic response of tall buildings was highlighted by the Commerce Court study of the 1970s (Section 9.2), when a building of a uniform rectangular cross-section experienced significant and measurable dynamic twist due to an eccentricity between the elastic and mass centres. Such a possibility had been overlooked in the original wind tunnel testing. Now, when considering accelerations at the top of tall building, the possibility of torsional motions increasing the perceptible motions at the periphery of the cross-section may need to be considered.

There are two mechanisms for producing dynamic torque and torsional motions in tall buildings:

- Mean torque and torsional excitation resulting from non-uniform pressure distributions, or from non-symmetric cross-sectional geometries, and
- Torsional response resulting from sway motions through coupled mode shapes and/or eccentricities between elastic (shear) and geometric centres.

The first aspect was studied by Isyumov and Poole (1983), Lythe and Surry (1990), and Cheung and Melbourne (1992). Torsional response of tall buildings has been investigated both computationally making use of experimentally obtained dynamic pressure or force data from wind tunnel models (Tallin and Ellingwood, 1985; Kareem, 1985), and exper-

imentally on aeroelastic models with torsional degrees of freedom (Xu *et al.*, 1992a; Beneke and Kwok, 1993; Zhang *et al.*, 1993).

A mean torque coefficient, \bar{C}_{M_z} , can be defined as:

$$\bar{C}_{M_z} = \frac{\bar{M}_z}{\frac{1}{2}\rho_a \bar{U}_h^2 b_{\max}^2 h} \tag{9.22}$$

where \bar{M}_z is the mean torque, b_{\max} is the maximum projected width of the cross-section, h is the height of the building.

Lythe and Surry (1990), from wind tunnel tests on sixty-two buildings, ranging from those with simple cross-sections to complex shapes, found an average value of \bar{C}_{M_z} , as defined above, of 0.085, with a standard deviation of 0.04. The highest values appear to be a function of the ratio of the minimum projected width, b_{\min} to the maximum projected width, b_{\max} , with a maximum value of \bar{C}_{M_z} approaching 0.2, when (b_{\min}/b_{\max}) is equal to around 0.45 (Figure 9.11 from Cheung and Melbourne, 1992). The highest value of \bar{C}_{M_z} for any section generally occurs when the mean wind direction is about 60–80 degrees from the normal to the widest building face.

Isyumov and Poole (1983) used simultaneous fluctuating pressures and pneumatic averaging (Section 7.5.2) on building models with a square or 2:1 rectangular cross-section in a wind tunnel, to determine the contribution to the fluctuating torque coefficient from various height levels on the buildings, and from the various building faces. The main contribution to the fluctuating torque on the square and rectangular section with the wind parallel to the long faces, came from pressures on the side faces, and could be predicted from the mean torque by quasi-steady assumptions (Section 4.6.2). On the other hand, for

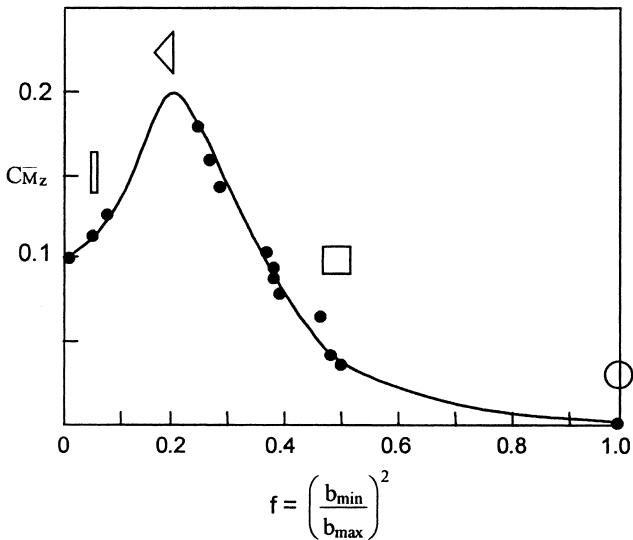


Figure 9.11 Mean torque coefficients on tall buildings of various cross sections (Cheung and Melbourne, 1992).

a mean wind direction parallel to the short walls of the rectangular cross-section, the main contribution was pressure fluctuations on the rear face, induced by vortex shedding.

A double peak in the torque spectra for the wind direction parallel to the long face of a 2:1 building has been attributed to buffetting by lateral turbulence, and by re-attaching flow on to the side faces (Xu *et al.*, 1992a). Measurements on an aeroelastic wind tunnel tall building model designed only to respond torsionally (Xu *et al.*, 1992a), indicated that aerodynamic damping effects (Section 5.5.1) for torsional motion of cross-section shapes characteristic of tall buildings are quite small in the range of design reduced velocities, in contrast to bridge decks. However at higher reduced velocities, high torsional dynamic response and significant negative aerodynamic damping has been found for a triangular cross-section (Beneke and Kwok, 1993).

A small amount of eccentricity can increase both the mean twist angle and dynamic torsional response. For example for a building with square cross-section, a shift of the elastic centre from the geometric and mass centre by 10% of the breadth of the cross-section, is sufficient to double the mean angle of twist and increase the dynamic twist by 40–50% (Zhang *et al.*, 1993).

9.8 Interference effects

High-rise buildings are most commonly clustered together in groups – as office buildings grouped together in a city-centre business district, or in multiple building apartment developments, for example. The question of aerodynamic interference effects from other buildings of similar size on the structural loading and response of tall buildings arises.

9.8.1 Upwind building

A single similar upwind building on a building with square cross-section and height/width (aspect) ratio of six produces increases of up to 30% in peak along-wind base moment, and 70% in cross-wind moment, at reduced velocities representative of design wind conditions in suburban approach terrain (Melbourne and Sharp, 1976). The maximum increases occur when the upwind building is two to three building widths to one side of a line taken upwind, and about eight building widths upstream. Contours of percentage increases in peak cross-wind loading for square-section buildings with an aspect ratio of 4, are shown in Figure 9.12. It can be seen that reductions, i.e. shielding, occurs when the upstream building is within four building heights upstream and ± 2 building heights to one side of the downstream building. The effect of increasing turbulence in the approach flow, i.e. increasing roughness lengths in the approach terrain, is to reduce the increases produced by interference.

The effect of increasing aspect ratio is to further increase the interference effects of upstream buildings, with increases of up to 80% being obtained, although this was for buildings with an atypical aspect ratio of 9, and in relatively low turbulence conditions. (Bailey and Kwok, 1985).

9.8.2 Downwind building

As shown in Figure 9.12, downwind buildings can also increase cross-wind loads on buildings if they are located in particular critical positions. In the case of the buildings of 4:1 aspect ratio of Figure 9.12, this is about one building width to the side, and two widths downwind.

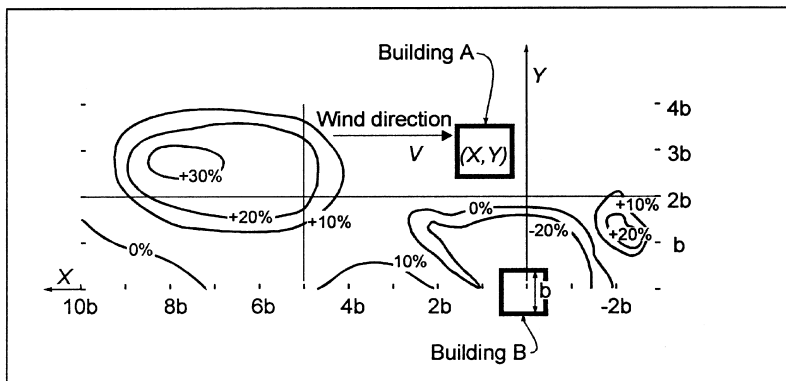


Figure 9.12 Percentage change in cross-wind response of a building (B) due to a similar building (A) at (X, Y) (Standards Australia, 1989).

More detailed reviews of interference effects on wind loads on tall buildings are given by Kwok (1995) and Khanduri *et al.* (1998). For a complex of tall buildings in the centre of large cities, wind tunnel model tests (Chapter 7) will usually be carried out, and these should reveal any significant interference effects on *new* buildings, such as those described in the previous paragraphs. Anticipated new construction should be included in the models when carrying out such tests. However, existing buildings may be subjected to unpredicted higher loads produced by new buildings of similar size at any time during their future life, and this should be considered by designers, when considering load factors.

9.9 Damping

The dynamic response of a tall building or other structure, to along-wind or cross-wind forces, depends on its ability to dissipate energy, known as ‘damping’. Structural damping is derived from energy dissipation mechanisms within the material of the structure itself (i.e. steel, concrete, etc.), or from friction at joints or from movement of partitions, etc. For some large structures constructed in the last twenty years, the structural damping alone has been insufficient to limit the resonant dynamic motions to acceptable levels for serviceability considerations, and auxiliary dampers have been added. Three types of auxiliary damping devices will be discussed in this chapter: viscoelastic dampers, tuned mass dampers (T.M.D.) and tuned liquid dampers (T.L.D.).

9.9.1 Structural damping

An extensive database of free vibration measurements from tall buildings in Japan has been collected (Tamura *et al.*, 2000). This database includes data on frequency as well as damping. More than 200 buildings were studied, although there is a shortage of values at larger heights – the tallest (steel encased) reinforced concrete building was about 170 m in height, and the highest steel-framed building was 280 m.

For reinforced concrete buildings, the Japanese study proposed the following empirical formula for the critical damping ratio in the first mode of vibration, for buildings less than 100 m in height, and for low-amplitude vibrations (drift ratio, (x_r/h) less than 2×10^{-5}).

$$\eta_1 \cong 0.014n_1 + 470\left(\frac{x_t}{h}\right) - 0.0018 \quad (9.23)$$

where n_1 is the first mode natural frequency, and x_t is the amplitude of vibration at the top of the building ($z=h$).

The corresponding relationship for steel-framed buildings is:

$$\eta_1 \cong 0.013n_1 + 400\left(\frac{x_t}{h}\right) + 0.0029 \quad (9.24)$$

The range of application for equation (9.24) is stated to be: $h < 200$ metres, and (x_t/h) less than 2×10^{-5} .

Equations (9.23) and (9.24) may be applied to tall buildings for serviceability limit states criteria (i.e. for the assessment of acceleration limits). Much higher values are applicable for the high amplitudes appropriate to strength (ultimate) limit states, but unfortunately little, or no, measured data are available.

9.9.2 Visco-elastic dampers

Visco-elastic dampers incorporate visco-elastic material which dissipates energy as heat through shear stresses in the material. A typical damper, as shown in Figure 9.13, consists of two visco-elastic layers bonded between three parallel plates (Mahmoodi, 1969). The force versus displacement characteristic of such a damper forms a hysteresis loop as shown in Figure 9.14. The enclosed area of the loop is a measure of the energy dissipated per cycle, and for a given damper, is dependent on the operating temperature (Mahmoodi and Keel, 1986) and heat transfer to the adjacent structure.

The World Trade Center buildings in New York City were the first major structures to utilise visco-elastic dampers (Mahmoodi, 1969). Approximately 10,000 dampers were

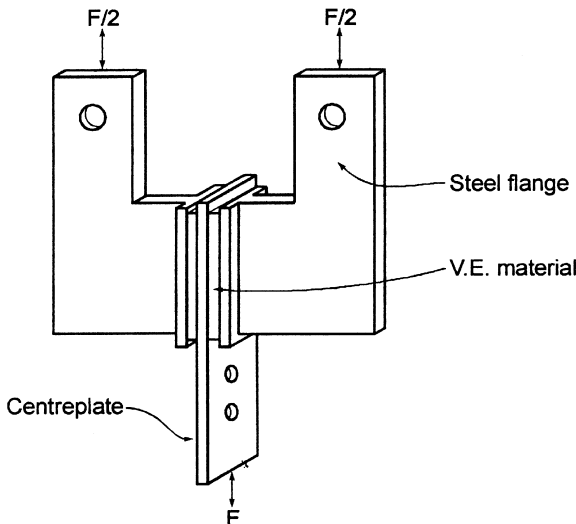


Figure 9.13 A viscoelastic damper (Mahmoodi, 1969).

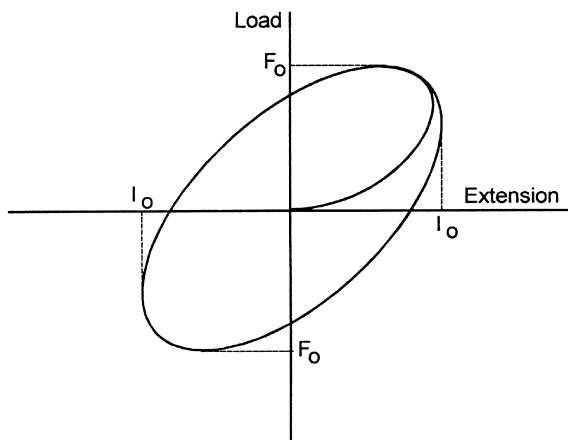


Figure 9.14 Hysteresis loop for viscoelastic damper (Mahmoodi, 1969).

installed in each 110-storey tower, with about 100 dampers at the ends of the floor trusses at each floor from the 7th to the 107th. More recently visco-elastic dampers have been installed in the 76-storey Columbia Seafirst Center Building, in Seattle, U.S.A. The dampers used in this building were significantly larger than those used at the World Trade Center, and only 260 were required to effectively reduce accelerations in the structure to acceptable levels (Skilling *et al.*, 1986; Keel and Mahmoodi, 1986).

A detailed review of the use of visco-elastic dampers in tall buildings has been given by Samali and Kwok (1995).

9.9.3 Tuned mass dampers

A relatively popular method of mitigating vibrations has been the *tuned mass damper* (T.M.D.) or *vibration absorber*. Vibration energy is absorbed through the motion of an auxiliary or secondary mass connected to the main system by viscous dampers. The characteristics of a vibrating system with T.M.D. can be investigated by studying the two-degree-of-freedom system shown in Figure 9.15 (e.g. den Hartog, 1956; Vickery and Davenport, 1970).

Tuned mass damper systems have successfully been installed in the Sydney Tower in Australia, the Citycorp Center, New York (275 m), the John Hancock Building, Boston, U.S.A. (60 storeys), and in the Chiba Port Tower in Japan (125 m). In the first and last of these, extensive full-scale measurements have been made to verify the effectiveness of the systems.

For the Sydney Tower, a 180-tonne doughnut-shaped water tank, located near the top of the Tower, and required by law for fire protection, was incorporated into the design of the T.M.D. The tank is 2.1 m deep and 2.1 m from inner to outer radius, weighs about 200 tonnes, and is suspended from the top radial members of the turret. Energy is dissipated in eight shock absorbers attached tangentially to the tank and anchored to the turret wall. A 40-tonne secondary damper is installed lower down on the tower to further increase the damping, particularly in the second mode of vibration (Vickery and Davenport, 1970; Kwok, 1984).

The system installed in the Citycorp Center Building, New York, (McNamara, 1977),

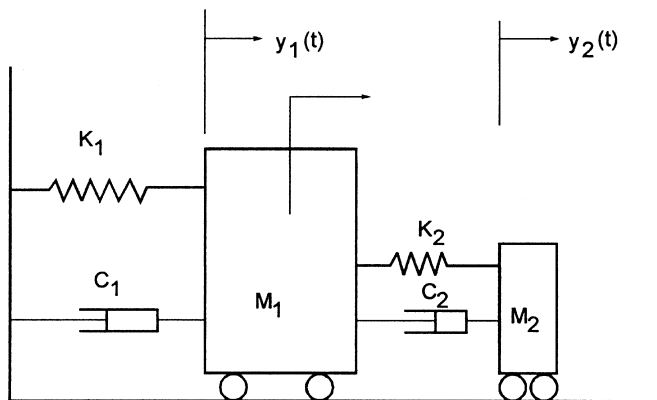


Figure 9.15 Two degree-of-freedom representation of a tuned mass damper.

consists of a 400-tonne concrete mass riding on a thin oil film. The damper stiffness is provided by pneumatic springs, whose rate can be adjusted to match the building frequency. The energy absorption is provided by pneumatic shock absorbers, as for the Sydney Tower. The building was extensively wind tunnel tested (Isyumov *et al.*, 1975). The aeroelastic model tests included the evaluation of the tuned mass damper. The T.M.D. was found to significantly reduce the wind-induced dynamic accelerations to acceptable levels. The effective damping of the model damper was found to be consistent with theoretical estimates of effective viscous damping based on the two-degree-of-freedom model (Vickery and Davenport, 1970).

T.M.D. systems similar to those in the Citycorp Building have been installed in both the John Hancock Building, Boston, and in the Chiba Port Tower. In the case of the latter structure, the system has been installed to mitigate vibrations due to both wind (typhoon) and earthquake. Adjustable coil springs are used to restrain the moving mass, which is supported on frames sliding on rails in two orthogonal directions.

The performance of tuned mass dampers in tall buildings and towers under wind loading has been reviewed by Kwok and Samali (1995).

9.9.4 Tuned liquid dampers

Tuned liquid dampers are relatively new devices in building and structures applications, although similar devices have been used in marine and aerospace applications for many years. They are similar in principle to the tuned mass damper, in that they provide a heavily damped auxiliary vibrating system attached to the main system. However the mass, stiffness and damping components of the auxiliary system are all provided by moving liquid. The stiffness is in fact gravitational; the energy absorption comes from mechanisms such as viscous boundary layers, turbulence or wave breaking, depending on the type of system. Two categories of T.L.D. will be discussed briefly here: *tuned sloshing dampers* (T.S.D.) and *tuned liquid column dampers* (T.L.C.D.).

The tuned sloshing damper type (Figure 9.16) relies on the motion of shallow liquid in a rigid container for absorbing and dissipating vibrational energy (Fujino *et al.*, 1988; Sun *et al.*, 1989). Devices of this type have already been installed in at least two structures in Japan (Fujii *et al.*, 1990) and on a television broadcasting tower in Australia.

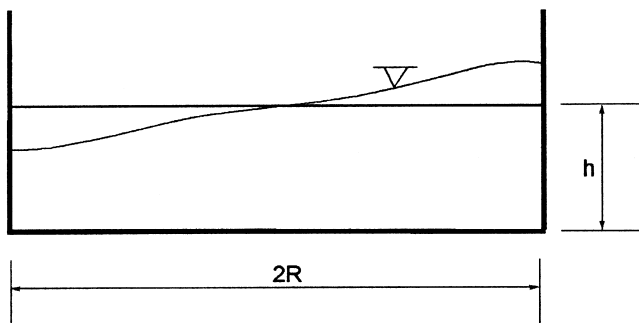


Figure 9.16 Tuned sloshing damper.

Although a very simple system in concept, the physical mechanisms behind this type of damper are in fact quite complicated. Parametric studies of dampers with circular containers were carried out by Fujino *et al.* (1988). Some of their conclusions can be summarised as follows:

- Wave breaking is dominant mechanism for energy dissipation but not the only one.
- The additional damping produced by the damper is highly dependent on the amplitude of vibration.
- At small to moderate amplitudes, the damping achieved is sensitive to the frequency of sloshing of liquid in the container. For dampers with circular containers, the fundamental sloshing frequency is given by equation (9.25).

$$n_s = (1/2\pi)\sqrt{[(1.84g/R)\tanh(1.84h/R)]} \quad (9.25)$$

where g is the acceleration due to gravity, h is the height of the liquid and R is the radius of the container, as shown in Figure 9.16. This formula is derived from linear potential theory of shallow waves.

- High viscosity sloshing liquid is not necessarily desirable at high amplitudes of vibration, as wave breaking is inhibited. However, at low amplitudes, at which energy is dissipated in the boundary layers on the bottom and side walls of the container, there is an optimum viscosity for maximum effectiveness (Sun *et al.*, 1989).
- Roughening the container bottom does not improve the effectiveness because it has little effect on wave breaking.

The above conclusions were based on a limited number of free vibration tests with only two diameters of container. Further investigations are required, including the optimal size of T.S.D. for a given mass of sloshing liquid. However the simplicity and low cost of this type of damper makes them very suitable for many types of structure.

Variations in the geometrical form are possible, for example Modi *et al.* (1990) has examined T.S.D.s with torus (doughnut)-shaped containers.

The ‘tuned liquid column damper’ (TLCD) damper (Figure 9.17) comprises an auxiliary vibrating system consisting of a column of liquid moving in a tube-like container. The restoring force is provided by gravity, and energy dissipation is achieved at orifices

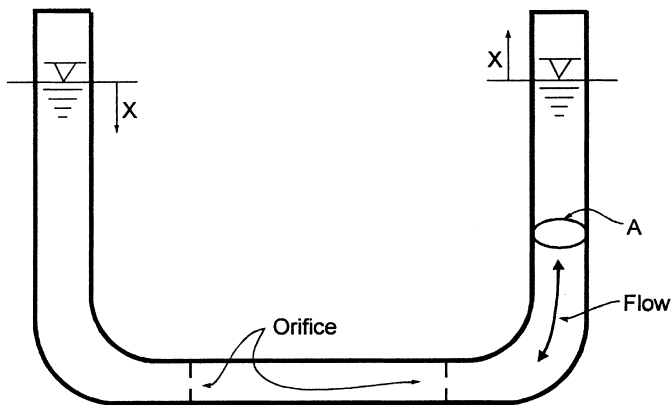


Figure 9.17 Tuned liquid column damper.

installed in the container (Sakai *et al.*, 1989; Hitchcock *et al.*, 1997a, 1997b). The same principle has been utilised in anti-rolling tanks used in ships.

The T.L.C.D., like the T.S.D., is simple and cheap to implement. Unlike the T.S.D., the theory of its operation is relatively simple and accurate. Sakai *et al.* (1989) has designed a T.L.C.D. system for the Citycorp Building, New York as a feasibility study; he found that the resulting damper was simpler, lighter and presumably cheaper than the T.M.D. system actually used in this building (Section 9.9.3). Xu *et al.* (1992b) have examined theoretically the along-wind response of tall, multi-degree-of-freedom structures, with T.M.D.s, T.L.C.D.s, and a hybrid damper – the Tuned Liquid Column Mass Damper (T.L.C.M.D.). They found that the T.M.D. and T.L.C.D., with the same amount of added mass, achieved similar response reductions. The T.L.C.M.D., in which the mass of the container, as well as the liquid, is used as part of the auxiliary vibrating system, is less effective when the liquid column frequency is tuned to the same frequency as the whole damper frequency (with the water assumed to remain still). The performance of the latter is improved when the liquid column frequency is set higher than the whole damper frequency.

The effectiveness of tuned liquid dampers in several tall structures in Japan has been reviewed by Tamura *et al.* (1995).

9.10 Case studies

Very many tall buildings have been studied in wind tunnels over several decades. These studies include the determination of the overall loading and response, cladding pressures, and other wind effects, such as environmental wind conditions at ground level. However, these studies are usually proprietary in nature, and not generally available. However, Willford (1985) has described a response study for the Hong Kong and Shanghai Bank Building, Hong Kong. A detailed wind engineering study for a building of intermediate height, including wind loading aspects, is presented by Surry *et al.* (1977). Relatively few tall buildings have been studied in full scale for wind loads, although many have been studied for their basic dynamic properties (e.g. Tamura *et al.*, 2000). Case studies of wind-induced accelerations on medium height buildings are described by Wyatt and Best (1984), and Snaebjornsson and Reed (1991).

9.11 Summary

This chapter has discussed various aspects of the design of tall buildings for wind loads. The general characteristics of wind pressures on tall buildings, and local cladding loads have been considered. The special response characteristics of glass have been discussed. The overall response of tall buildings in along-wind and cross-wind directions, and in twist (torsion) has been covered. Aerodynamic interference effects, and the application of auxiliary damping systems to mitigate wind-induced vibration have been discussed.

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